Shatt Al-Arab University College Department Civil Eng. Fluid Mechanics-I **SECOND YEAR** Dr. Jasim Al-Battat

The Syllabus

- **1.** Introductory Concepts of Fluid Mechanics
- **2.** Properties of a Fluid
- 3. Fluid Static
- 4. Fluid Dynamics

References

- 1. Fluid Mechanics by Streeter and Wylie.
- 2. Fluid Mechanics for Engineers by Albertson, Barton, and Simons.
- 3. Fluid Mechanics by Hydraulics (Schaum's Series) by Griles
- ميكانيك الموائع دز نزار علي سبتي 4.
- مبادئ ميكانيك الموائع د. جميل المالئكة 5.
- 6. Fluid Mechanics with Engineering Applications by Daugherty, Franzini, and Finnemore.
- 7. Elementary Fluid Mechanics by Vennard and Street.

Chapter One

Introductory Concepts of Fluid Mechanics

1.1. The Concept of a Fluid and Fluid Mechanics

Mechanics is the oldest physical science that deals with both stationary and moving bodies under the influence of forces. The branch of mechanics that deals with bodies at rest is called statics, while the branch that deals with bodies in motion is called dynamics. The subcategory fluid mechanics is defined as the science that deals with the behavior of fluids at rest (fluid statics) or in motion (fluid dynamics), and the interaction of fluids with solids or other fluids at the boundaries. Fluid mechanics is also referred to as fluid dynamics by considering fluids at rest as a special case of motion with zero velocity. Fluid is a substance that deforms continuously when subjected to shear stress, no matter how small that shear stress may be. Fluids may be either liquids or gases. Solids, as compared to fluids, cannot be deformed permanently (plastic deformation) unless a certain value of shear stress (called the yield stress) is exerted on it. Figure 1.1 illustrates a solid block resting on a rigid plane and stressed by its own weight. The solid sags into a static deflection, shown as a highly exaggerated dashed line, resisting shear without flow. A free-body diagram of element A on the side of the block shows that there is shear in the block along a plane cut at an angle θ through A. Since the block sides are unsupported, element A has zero stress on the left and right sides and compression stress $\sigma = -p$ on the top and bottom. Mohr's circle does not reduce to a point, and there is nonzero shear stress in the block.

Prior to fluid mechanics, statics, and dynamics, was taken, involve solid mechanics. Mechanics is the field of science focused on the motion of material bodies. Mechanics involves force, energy, motion, deformation, and material properties. When mechanics applies to material bodies in the solid phase, the discipline is called solid mechanics. When the material body is in the gas or liquid phase, the discipline is called fluid mechanics. In contrast to a solid, a fluid is a substance whose molecules move freely past each other. More specifically, a fluid is a substance that will continuously deform [that is, flow under

the action of a shear stress]. Alternatively, a solid will deform under the action of a shear stress but will not flow like a fluid. Both liquids and gases are classified as fluids.

This lecture notes introduces fluid mechanics by describing gases, liquids, and the continuum assumption. This lecture notes also presents an approach for using units and primary dimensions in fluid mechanics calculations.

1.2 Liquids and Gases

Liquids and gases differ because of forces between the molecules. As shown in the figure 1.1, a liquid will take the shape of a container whereas a gas will expand to fill a closed container. The behavior of the liquid is produced by strong attractive force between the molecules. This strong attractive force also explains why the density of a liquid is much higher than the density of gas. A *gas* is a phase of material in which molecules are widely spaced, molecules move about freely, and forces between molecules are minuscule, except during collisions. Alternatively, a *liquid* is a phase of material in which molecules are closely spaced, molecules move about freely, and there are strong attractive forces between molecules.

1.3 Application Areas of Fluid Mechanics

Why are we studying fluid mechanics on a Civil Engineering course? The provision of adequate water services such as the supply of potable water, drainage, sewerage is essential for the development of industrial society. It is these services which civil engineers provide. Fluid mechanics is involved in nearly all areas of Civil Engineering either directly or indirectly. Some examples of direct involvement are those where we are concerned with manipulating the fluid:

- Sea and river (flood) defenses:
- Water distribution / sewerage (sanitation) networks;
- Hydraulic design of water/sewage treatment works;
- Dams;
- Irrigation;
- Pumps and Turbines;
- Water retaining structures.

And some examples where the primary object is construction - yet analysis of the fluid mechanics is essential:

- Flow of air around buildings;
- Bridge piers in rivers;
- Ground-water flow.

1.4 Dimensions and Units,

A dimension is the measure by which a physical variable is expressed quantitatively. A unit is a particular way of attaching a number to the quantitative dimension. In fluid mechanics there are only four primary dimensions from which all other dimensions can be derived: mass, length, time, and temperature. These dimensions and their units in both systems are given in Table 1.1. Note that the kelvin unit uses no degree symbol. The braces around a symbol like [M] mean "the dimension" of mass. All other variables in fluid mechanics can be expressed in terms of [M], [L], [T], and [Θ]. For example, acceleration has the dimensions $[LT^2]$.

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Table 1.1: Primary Dimensions in SI and BG Systems.

A list of some important secondary variables in fluid mechanics, with dimensions derived as combinations of the four primary dimensions, is given in Table 1.2.

Table 1.2: Secondary Dimensions in Fluid Mechanics.

Secondary dimension	SI unit	BG unit	Conversion factor
Area $\{L^2\}$	m ²	ft ²	$1 m2 = 10.764 ft2$
Volume $\{L^3\}$	m ³	ft ³	$1 m3 = 35.315 ft3$
Velocity $\{LT^{-1}\}$	m/s	ft/s	$1 \text{ ft/s} = 0.3048 \text{ m/s}$
Acceleration $\{LT^{-2}\}$	m/s ²	ft/s^2	$1 \text{ ft/s}^2 = 0.3048 \text{ m/s}^2$
Pressure or stress			
$\{ML^{-1}T^{-2}\}$	$Pa = N/m2$	1 ¹	1 lbf/ft ² = 47.88 Pa
Angular velocity $\{T^{-1}\}\$	s^{-1}	s^{-1}	$1 s^{-1} = 1 s^{-1}$
Energy, heat, work			
$\{ML^{2}T^{-2}\}$	$J = N \cdot m$	ft·lbf	$1 \text{ ft} \cdot \text{lbf} = 1.3558 \text{ J}$
Power $\{ML^{2}T^{-3}\}$	$W = J/s$	ft·lbf/s	1 ft \cdot lbf/s = 1.3558 W
Density $\{ML^{-3}\}$	kg/m ³	slugs/ ft^3	1 slug/ft ³ = 515.4 kg/m ³
Viscosity $\{ML^{-1}T^{-1}\}$	$kg/(m \cdot s)$	slugs/ $(ft \cdot s)$	$1 \text{ slug/(ft} \cdot \text{s}) = 47.88 \text{ kg/(m} \cdot \text{s})$
Specific heat $\{L^2T^{-2}\Theta^{-1}\}$	$m^2/(s^2 \cdot K)$	$ft^2/(s^2 \cdot \text{°R})$	$1 \text{ m}^2\text{/s}^2 \cdot \text{K}$ = 5.980 ft ² /(s ² · °R)

Example 1:

A body weighs 1000 lbf when exposed to a standard earth gravity $g = 32.174$ ft/s². (a) What is its mass in kg ?

 (b) What will the weight of this body be in N if it is exposed to the moon's standard acceleration $g_{\text{moon}} = 1.62 \text{ m/s}^2$?

 (c) How fast will the body accelerate if a net force of 400 lbf is applied to it on the moon or on the earth?

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Solution:

Part (a) We can express Eq. (1) dimensionally, using braces by entering the dimensions of each term from Table 1.2:

$$
\{ML^{-1}T^{-2}\} = \{ML^{-1}T^{-2}\} + \{ML^{-3}\}\{L^{2}T^{-2}\} + \{ML^{-3}\}\{LT^{-2}\}\{L\}
$$

=
$$
\{ML^{-1}T^{-2}\}
$$
 for all terms

Part (b) Enter the SI units for each quantity from Table 1.2:

The right-hand side looks bad until we remember from Eq. (1.3) that 1 kg = 1 N · s²/m.

$$
\{kg/(m \cdot s^2)\} = \frac{\{N \cdot s^2/m\}}{\{m \cdot s^2\}} = \{N/m^2\}
$$
 Ans. (b)

Thus all terms in Bernoulli's equation will have units of pascals, or newtons per square meter, when SI units are used. No conversion factors are needed, which is true of all theoretical equations in fluid mechanics.

Part (c) Introducing BG units for each term, we have

> ${lbf/ft²} = {lbf/ft²} + {slugs/ft³} {ft²/s²} + {slugs/ft³} {ft/s²} {ft}$ = { lbf/ft^2 } + { $slugs/(ft \cdot s^2)$ }

But, from Eq. (1.3), 1 slug = 1 lbf \cdot s²/ft, so that

$$
{\{slugs/(ft \cdot s^2)\}} = {\frac{\{lbf \cdot s^2/ft\}}{\{ft \cdot s^2\}}} = {\{lbf/ft^2\}}
$$
 Ans. (c)

1 (Ibf) = 4.45 (N)

1 (ft) = 0.305 (m)

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Chapter Two

Properties of a Fluid

By Newton's second law,
$$
F=wa \Rightarrow W=m.3
$$

\n
$$
\therefore Y = \frac{w.9}{4} \Rightarrow Y = P3
$$
\n
$$
\Rightarrow Y = \frac{w.9}{4} \Rightarrow Y = P3
$$
\n
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\Rightarrow Y = \frac{w.9}{4} \Rightarrow Y = P3
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\Rightarrow Y = \frac{w.9}{4} \Rightarrow Y = P3
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$$
\Rightarrow \frac{w.9}{4} \Rightarrow \frac{w.9
$$

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Liquids are ordinary considered incompressible fluid since the change in volume (or density per unit mass) is so small of can be neglected = CE - constant). When (E) change, this means that the fluid is compressible as in air (gases in general). 3 - Viscosity. The viseosity of a fluid is a measure of its resistance to shear on any deformation. The friction forces in fluid flow result from cohesion & momentum interchange between molecules in the fluid. As the temperature increases, the viscosities of all liquids decrease, while the viscosities of all gases increase. This is because the force of cohesion, which diminishes with temperature, predominates with liquids, while with gases the predominating factor is the interchange of molecules between layers of different velocities. - M (mu) (dynamic viscosity) $(\frac{N\cdot s}{m^2}$ = Pascal.s. M_{water} = 1.005 $\times 10^{-3}$ pa. s. $a + 22$ M_2 _{di} C_2 = 1.8 * 10 $Pa. s.$
atzo - $N(nu)$ (Kinematic viscosity) (m²) $N = \frac{N}{P}$ $N = \frac{\frac{N.5}{M^2}}{N.5} = \frac{N.5}{M} \cdot \frac{M}{N.5^{2}}$ \therefore N = $\frac{m^2}{s}$.

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4-Vapor Pressure: All liquids fend to evaporate or vaporize, which they do by projecting molecules into the space above their surface. Molecular activity increases with temperature, & hence vapor pressure increases with temperature. The phenomena of vaporization & boiling are differentiated as follows:

* Vaporization if (Vapor pressure < pressure above aliquid surface at atemperature)

* Boiling if (Vapor pressure = pressure above a liquid surface at a temperature) to the Pressure a bout 4 4 p vapor pressure

(Sigma)
5 - Surface Tension (J) (Euglin = M)

Liquids have cahesion of adhesion, both of which are forms of molecular attraction. Cohesion enables a liquid to resist tensile stress, while adhesion enables it to adhere to another body. The attraction between molecules forms an imaginary film capable of resisting tension at the interface between two immiscible liquids or at the interface between a liquid & a gas. The liquid property that creates this capability is known as surface Tension.

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From
$$
-q.0
$$
: $\sqrt{\cos\theta}.\vec{x}d=\sqrt{\frac{x}{4}}d^{2}h$
\n $\therefore h = \frac{4\alpha cos\theta}{6}d$
\nwhere, $h = Capitala\theta + k^2\theta + k^3\theta + k^4\theta$
\n $\alpha = 5x+4\theta + k^3\theta + (N/a^3) = 93$
\n $\beta = 1:9x+3$ density $(K3\ln^3)$
\n $\beta = 9\pi\sqrt{13} d$ density $(K3\ln^3)$
\n $\beta = 9\pi\sqrt{13} d$ acceleration (where?)
\n $\beta = 1:3\pi/3 d$ acceleration (where?)
\n $\beta = 0$ for water
\n $\beta = 130$ If mercury (H3)
\n $\beta = 130$ If mercury (H3)
\n $\beta = 130$ If mercury (H4)
\n $\beta = 130$ If mercury (H5)
\n $\beta = 130$ If water
\n $\beta = 130$

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Solution:
$$
\frac{a}{\pm}
$$
 Since the glass is clean = $\frac{a}{\pm}$ = 0
\nh = $\frac{4\sqrt{cos\theta}}{\sqrt{d}}$ = $\frac{4\sqrt{cos\theta}}{8d}$ = $\frac{4\sqrt{cos\theta}}{8d}$
\nh = $\frac{4*0.0736}{9810*4*10^{3}}$ = 0.0075m = 7.5mm rise
\n+ unit of the change 130 km/h³ = 0.00246 m
\n $\frac{4*0.51*cos130}{13.6*8810*4*10} = -2.46$ mm
\ndepression of

7. Temperature Dependency

The effect of temperature on viscosity is different for liquids and gases. The viscosity of liquids decreases as the temperature increases, whereas the viscosity of gases increases with increasing temperature. To understand the mechanisms responsible for an increase in temperature that causes a decrease in viscosity in a liquid, it is helpful to rely on an approximate theory that has been developed to explained the observed trends (1). The molecules in a liquid form a structure with "holes" where there are no molecules, as shown in Fig. 2.2. Even when the liquid is at rest, the molecules are in constant motion, but confined to cells. The cell structure is caused by attractive forces between the molecules. The cells may be thought of as energy barriers. When the liquid is subjected to a rate of strain and thus caused to move, as shown in Fig. 2.2, there is a shear stress, τ , imposed by one layer on another in the fluid. This force/area assists a molecule in overcoming the energy barrier, and it can move into the next hole. The magnitude of these energy barriers is related to viscosity, or resistance to shear deformation. At a higher temperature the size

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of the energy barrier is smaller, and it is easier for molecules to make the jump, so that the net effect is less resistance to deformation under shear. Thus, an increase in temperature causes a decrease in viscosity for liquids. An equation for the variation of liquid viscosity with temperature is

$$
\mu = Ce^{b/T} \tag{2.9}
$$

where *C* and *b* are empirical constants that require viscosity data at two temperatures for evaluation.

Fig. 2.2

As compared to liquids, gases do not have zones to which molecules are confined by intermolecular bonding. Gas molecules are always undergoing random motion. If this random motion of molecules is superimposed upon two layers of gas, where the top layer is moving faster than the bottom layer, periodically a gas molecule will randomly move from one layer to the other. As the gas temperature increases, more of the molecules will be making random jumps. Just highly mobile gas molecules have momentum, which must be resisted by the layer to which the molecules jump. Therefore, as the temperature increases, the viscosity, or resistance to shear, also increases.

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Example 4

CALCULATING VISCOSITY OF LIQUID AS A FUNCTION OF TEMPERATURE

The dynamic viscosity of water at 20 $^{\circ}$ C is 1.00×10^{-3} N · s/m² and the viscosity at 40 $^{\circ}$ C is

 6.53×10^{-4} N·s/m². Estimate the viscosity at 30°C. Viscosity of water is

specified at two temperatures. Find The viscosity at 30°C by interpolation.

a) Water at 20 $^{\circ}$ C, $\mu = 1.00 \times 10^{-3}$ N · s/m².

b) Water at 40°C, $\mu = 6.53 \times 10^{-4} \text{ N} \cdot \text{s/m}^2$

Solution

- 1. Logarithm of Eq. (2.9)
- $\ln \mu = \ln C + b/T$
- 2. Interpolation

 $-6.908 = \ln C + 0.00341b$

 $-7.334 = \ln C + 0.00319b$

3. Solution for and *b*

 $ln C = -13.51$ $b = 1936$ (K)

 $C = e^{-13.51} = 1.357 \times 10^{-6}$

4. Substitution back in exponential equation

 $\mu = 1.357 \times 10^{-6} e^{1936/T}$

At 30°C

 $\mu = 8.08 \times 10^{-4} \text{ N} \cdot \text{s} / \text{m}^2$

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T = Torque in the torsional spring.
\n
$$
w =
$$
 constant angular velocity (rad. |s.) = $\frac{2\pi N}{60}$
\nN = angular velocity (r.p.m.) (revolutionger minute).
\n
\n $\frac{1}{\pi}$ There are two types of fluids depending on the existing of W
\n $\frac{1}{\pi}$ Read (Non-Viscaus) $\Rightarrow W=0$.
\n
\nFinally
\n $\frac{1}{\pi}$ Real (Viscous) $\Rightarrow W \neq 0$.
\n $\frac{1}{\pi}$ Nenn-Newtonian (w-Uories)
\n $\frac{1}{\pi}$ Nenn-Newtonian (w-Uories)
\n $\frac{1}{\pi}$ Auchy-
\n $\frac{1}{\pi}$ Auchy-
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\n $\frac{1}{\pi}$ Auchy-
\n $\frac{1}{\pi}$ Aholy-
\n $\frac{1}{\pi}$ Ahol

w = w
is uaried

Ex1: Find the required force (F) to pull within plate at the middle point between two large plates, which the distance between them (0.02m). The fluid between plates has
$$
(\mu = 0.862 \text{ Pa} \cdot \text{s})
$$
. The fluid between plates has $(\mu = 0.862 \text{ Pa} \cdot \text{s})$. The surface area of the thin plate is (0.465 m) for each face. The constant velocity of the thin plate is (0.152 m/s) .

\nNote: Assume linear velocity distribution.

Solution:
\nFor constant velocity of
\n
$$
\pi r
$$
 and πr is
\n πr and πr is
\n $r = m \cdot a \Rightarrow F = 0$.
\n $\therefore \Sigma F_{X=0}$.
\n $F = F_1 + F_2$
\n $F_1 = F_2 = n \cdot A$
\n $\Rightarrow F = 2F_1 = 2n \cdot A$
\n $\Rightarrow F = 2F_1 = 2n \cdot A$
\n $\Rightarrow F = 2F_1 = 2n \cdot A$
\n $\therefore F = 0.862 (15.2) = 13.1 Pa$.
\n $\therefore F = 2(13.1) * 0.465 = 12.183 N$

Ex.2: A long circular rod of (70 nm) diameter slides concern-
\nvically in (150 nm) long fixed tube, 3 of
\n(70.5 nm) interval diameter. The annular space between
\nthe real 4 the tube is filled with oil of viscosity
\n(0.193 Pa.s.). What force is required to slide the rod
\nthrough the tube with a velocity of (1m/s.)?

\nNote: Assume the velocity of the velocity of the velocity of the velocity.

\nUse: Assume the velocity of the velocity of the velocity of the velocity.

\nSolution:

\n
$$
\frac{100}{100} = \
$$

Ex.3: Water flows in a long pipe of dia. (0.305m). The velocity profile has a parabolic shape, (v=10y-32.8g) where (y) is the distance measured from pipe well toward the center. Find the equation of shear stress distribution, then calculate the shear stress at the wall 4 at the center of the pipe. $\begin{pmatrix} w & 1 & 1 \\ 0 & 1 & 1 \end{pmatrix}$ = 1.307 \star (0) $\begin{pmatrix} 3 & 0 & 0 \\ 0 & 1 & 1 \end{pmatrix}$

A cylinder (50 mm) in radius of (0.6m) in length, $Ex.4$ rotates coaxially inside a fixed cylinder of radius(56 mm), as shown in figure below. Liquid of ($\mu = 1.48$ Pa.s.) fills the space between the two cylinders & between the inner cylinder of the base. Calculate the torque required to rotate the inner cylinder at a constant angular velocity of (201. p.m). (Take end effects into consideration). 15^{ω} liquid 0.6 $Sol: T = T_1 + T_2$ where, $T_1 =$ torque exerted on the velocity wall of the inner cylinder. T2 = torque exerted on the x Soum base of the inner cylinder $T = F_1 \cdot r = F_1(\text{0.05})$ $F_1 = F_1 \cdot A_1 = \mu \frac{du}{dy} A_1$ $\frac{dv}{dy} = \frac{v_2 - v_1}{v_2 - v_1} = \frac{v_2}{0.006} = \frac{w \cdot r}{0.006} = \frac{w \cdot (0.05)}{0.006} = 8.384 \text{ W}$ $W = \frac{2 \times N}{60} = \frac{2 \times (20)}{60} = 2.09 \text{ rad.}$ 22

$$
\frac{3v}{dy} = 8.334 (2.09) = 17.425
$$
\n
$$
A_{1} = \text{surface area} + \text{inner cylinder} = 2\pi (0.05.)(0.6)
$$
\n
$$
= 0.188 \text{ m}^2
$$
\n
$$
\therefore F_{1} = 1.48 (17.42)(0.188) = 4.85 \text{ N}
$$
\n
$$
= 0.188 \text{ m}^2
$$
\n
$$
\therefore T_{1} = 4.85 (0.05) = 0.243 \text{ N} \cdot \text{m}
$$
\n
$$
T_{2} = F_{2} \cdot \text{m}
$$
\n
$$
F_{2} = P_{2} \cdot \text{m}
$$
\n
$$
F_{2} = \frac{W \cdot \text{m}}{32.31} = \frac{W \cdot \text{m}}{0.006}
$$
\n
$$
A_{2} = 2\pi \text{m}^2
$$
\n
$$
\therefore T_{2} = \frac{W \cdot \text{m}}{0.006} (2\pi \cdot \text{m}^2)
$$
\n
$$
= 0.05
$$
\n
$$
\therefore T_{2} = \frac{2\pi (1.48)(2.09)}{0.006} = 0.248 \text{ N} \cdot \text{m}
$$
\n
$$
T_{2} = 5.06 \times 10^{3} \text{ N} \cdot \text{m}
$$
\n
$$
\therefore T_{2} = 0.243 + 5.06 \times 10^{3} = 0.248 \text{ N} \cdot \text{m}
$$

Equation of state for Perfect Bias
\nThe perfect gas is defined as a substance that satisfies
\nthe perfect gas law g has a constant specific heats. The
\nequation of state for perfect gas is;
\n
$$
Pu = \frac{P}{P} = RT
$$
\nwhere;
\n
$$
P = absolute pressure (N/m^2 \leq R.)
$$
\nwhere;
\n
$$
P = absolute pressure (N/m^2 \leq R.)
$$
\nwhere;
\n
$$
P = absolute pressure (N/m^2 \leq R.)
$$
\n
$$
V = speed; i.2
$$
\n
$$
V = speed; i.2
$$
\n
$$
V = p = \sqrt{N}
$$
\nwhere,
$$
V = 1
$$
 and
$$
V = 1
$$
 and <math display="block</p>

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Op is defined as the number of heat units added per unit mass to raise the temperature of the gas one degree when the pressure is held constant.

The perfect gas must be carefully distinguished from ideal fluid. An ideal fluid is frictionless (N=0.) 4 incompress-: ble (E-const.). The perfect gas has viscosity & can there. fore develop shear stresses, of it is compressible according to the equation of state described above.

For Perfect gas, it is found;

 $Pv^n = P_1v_1^n = constant$

where; n = any non-negative value depending upon the process to which the gas is subjected.

Ex1. How many kilograms of Carbon monoxide gas at

\n2000 f (1001) if
$$
R = 297 \frac{J}{kg \cdot k}
$$

\nSolution.

\nmass = ? $\Rightarrow P = V$ (where $P = \frac{m}{H} \Rightarrow m = P + J$

\nsince $P = PRT$

\n $\therefore P = \frac{m}{RT} = \frac{200 \times 10}{297 (20 + 273)} = 2.298 \times 100 \times 10^3$

\nSince $P = PRT$

\n $\therefore P = \frac{m}{RT} = \frac{200 \times 10}{297 (20 + 273)} = 2.298 \times 100 \times 10^3$

\nSince $P = \frac{m}{H} \Rightarrow m = P + J = 2.298 \times 100 \times 10^3$

\n $\therefore m = 0.23 \text{ kg}$

\nEx2. A container holds (1 kg) air is added and the final temperature is 1100, determine the final absolute pressure. $(R_{air}) = 287 \frac{J}{kg \cdot k}$

\nSolution:

\nSolution:

\nSolution:

\nSince, $P = PRT$

\nProof:

\nTime = $\frac{15}{300}$

\nTime = $\frac{12.5}{9.16}$

\nTime = $\frac{12.5}{9.16$

$$
P_{final} = \frac{2.5}{\pi} (287)(383) - 0
$$

\n
$$
46 \text{ find } t; \text{ since } P_{in} = P_{in} \text{ RT}_{init}.
$$

\n
$$
9 * 10 = \frac{1}{\pi} (287)(30 + 273)
$$

\n
$$
9 * 10 = \frac{1}{\pi} (287)(30 + 273)
$$

\n
$$
1. \pm 9.66 * 10^{-3} \text{ m}
$$

\n
$$
F_{final} = \frac{2.5}{9.66 * 10^{-3}} (287)(383)
$$

\n
$$
F_{final} = 28.45 * 10^{-6} Pa.
$$

\n
$$
= 28.45 MPa.
$$

\n
$$
= 2.8.45 MPa.
$$

\n
$$
\frac{m-1}{T_1} = (\frac{p_2}{p_1})^{\frac{m-1}{n}}
$$

Chapter Three Fluid Static

This unit begins mechanics of fluids in depth by introducing many concepts related to pressure and by describing how to calculate forces associated with distributions of pressure. This chapter is restricted to fluids that are in hydrostatic equilibrium.

Pressure is defined as the ratio of normal force to area at a point. The pressure often varies from point to point. For example, pressure acting on the water tank wall will vary at different locations on the wall. Pressure is a scalar quantity; that is, it has magnitude only. Units for pressure give a ratio of force to area. Newtons per square meter of area, or Pascal (Pa), is the SI unit. The USC units include psi, which is pounds-force per square inch, and psf, which is poundsforce per square foot.

3.1 Absolute Pressure, Gage Pressure, and Vacuum Pressure

Absolute pressure is referenced to regions such as outer space, where the pressure is essentially zero because the region is devoid of gas. The pressure in a perfect vacuum is called absolute zero, and pressure measured relative to this zero pressure is termed *absolute pressure.* When pressure is measured relative to prevailing local atmospheric pressure, the pressure value is called *gage pressure.* For example, when a tire pressure gage gives a value of 300 kPa (44 psi), this means that the absolute pressure in the tire is 300 kPa greater than local atmospheric pressure. To convert gage pressure to absolute pressure, add the local atmospheric pressure. When pressure is less than atmospheric, the pressure can be described using vacuum pressure. *Vacuum pressure* is defined as the difference between atmospheric pressure and actual pressure.

Figure 3.1 provides a visual description of the three pressure scales. Gage, absolute, and vacuum pressure can be related using equations labeled as the "pressure equations."

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3.2 Pressure Variation with Elevation

3.2.1 Hydrostatic Differential Equation

The hydrostatic differential equation is derived by applying force equilibrium to a static body of fluid. To begin the derivation, isolate a small cylindrical body, and then sketch a free-body diagram (FBD) as shown in Fig. 3.2. The cylindrical body

is oriented so that its longitudinal axis is parallel to an arbitrary l direction. The body is Δl long, ΔA in cross-sectional area, and inclined at an angle α with the horizontal.

Fig. 3.2

Apply force equilibrium in the l direction:

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$$
\sum F_{\ell} = 0
$$

$$
F_{\text{Pressure}} - F_{\text{Weight}} = 0
$$

$$
p\Delta A - (p + \Delta p)\Delta A - \gamma \Delta A \Delta \ell \sin \alpha = 0
$$

Simplify and divide by the volume of the body $\Delta\ell\Delta A$ to give $\frac{\Delta p}{\Delta\ell} = -\gamma \sin \alpha$

 $\sin \alpha = \frac{\Delta z}{\Delta \ell}$ From Fig. 3.2, the sine of the angle is given by Combining the previous two equations and letting approach zero gives

$$
\lim_{\Delta z \to 0} \frac{\Delta p}{\Delta z} = -\gamma
$$

The final result is $\frac{dp}{dz} = -\gamma$ (hydrostatic differential equation) (3.1)

Equation (3.1) means that changes in pressure correspond to changes in elevation.

3.2.2 Hydrostatic Equation

The hydrostatic equation is used to predict pressure variation in a fluid with constant specific weight (constant density). Integrating Eq. (3.1) will give

$$
p + \gamma z = p_z = \text{constant} \tag{3.2}
$$

Where, *z*: is elevation, (vertical distance above a fixed reference point - datum),

pz; *piezometric pressure*.

Dividing Eq. (3.2) by γ gives

$$
\frac{p_z}{\gamma} = \left(\frac{p}{\gamma} + z\right) = h = \text{constant} \tag{3.3}
$$

Where, *h* : *piezometric head*. Since *h* is constant in Eq. (3.3),

 $\frac{p_1}{\gamma}$ + z₁ = $\frac{p_2}{\gamma}$ + z₂ $(3.4a)$

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Multiplying Eq. $(3.4a)$ by γ gives

$$
p_1 + \gamma z_1 = p_2 + \gamma z_2 \tag{3.4b}
$$

In Eq. (3.4b), letting $\Delta p = p_2 - p_1$ and $\Delta z = z_2 - z_1$ letting gives

$$
\Delta p = -\gamma \Delta z \tag{3.4c}
$$

The hydrostatic equation is given by either Eq. (3.4a), (3.4b), or (3. c). These three equations are equivalent because any one of the equations can be used to derive the other two. Piezometric pressure and head are related by

$$
p_z = h\gamma \tag{3.5}
$$

When hydrostatic equilibrium prevails in a body of fluid of constant density, then *h* will be constant at all locations.

3.3 Pressure Measurements

Four scientific instruments for measuring pressure: the barometer, Bourdontube gage, piezometer and manometer, transducer will described.

Measurement Devices asurement Devices
a-Tube gauge - piezometer (Column) 7 used for
a-Tube gauge - Manometer

b- Bourdon gauge.

a - Tube Gauge
1 - Piezometer (Column): it is a simple device which using for moderate (the) pressures of liquids. It consists of a tube in which the liquid can freely rise without overflowing.

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PA = Patw. + 8 h

\nsince
$$
Patm = 0
$$
. (Zero gauge)

\n1. $7A = 8h$

\n2 - Manometer

\ni. 5° in the Manometer

\nii. D : Here, and 2 is the same value of $ln(10000)$.

\niii. D : Here, and 2 is the same value of $ln(10000)$.

\ni. 5° in the image.

\nii. D : Here, and 2 is the same value of $ln(10000)$.

\niii. D : Here, and 2 is the same value of $ln(10000)$.

\niv. 2 is the same value of $ln(10000)$.

\niv. 2 is the same value of $ln(10000)$.

\niv. 2 is the same value of $ln(10000)$.

\niv. 2 is the same value of $ln(10000)$.

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\niv. 2 is the same value of $ln(10000)$.

\niv. 2 is the same value of $ln(10000)$.

\niv. 2 is the same value of $ln(10000)$.

\niv. 2 is the same value

 $P_c = P_D$
 $P_A + \delta_1 h_1 = P_{a_1} + \delta_2 h_2$ \therefore $P_A = \delta_2 h_2 - \delta_1 h_1$ + ve press.

Shatt Al-Arab University College/Department Civil Eng. Fluid Mechanics-I /second Year Dr. Jasim Mohsin IV-Inverted Differential Manometer. It is an inverted U_tube device used to Ω measure small pressure difference between two connected points. $P_c = P_D$ hι $PA-S, h_1 = PB - Y_2 h_2 - Y_1$ $- P_A - P_B = \delta_1 h_1 - \delta_2 h_2 - \delta_1 h_1$

b. Bourdon-Tube Gage

A *Bourdon-tube* gage measures pressure by sensing the deflection of a coiled tube. The tube has an elliptical cross section and is bent into a circular arc, as shown in Fig. 3.4*.*When atmospheric pressure (zero gage pressure) prevails, the tube is undeflected. When pressure is applied to the gage, the curved tube tends to straighten, thereby actuating the pointer to read a positive gage pressure.

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P. **CRAC**

EXECUTE:
$$
1.36 \text{ kg}
$$
 = 2.5
\n**EXECUTE:** $P_{E} = P_{F}$
\n $P_{A} = .3 \text{ m/s}$ (0.6) = P_{n-m}
\n $P_{B} = -9810 \text{ m.o.6}$
\n $P_{B} = -9810 \text{ m.o.6}$
\n $P_{B} = -9810 \text{ m.o.6}$
\n $P_{B} = -5886 \text{ N/m}^2$
\n $P_{B} = 5886 \text{ Pa}$ (Vacuum)
\n $P_{B} = 1905 \text{ N/m}^2$
\n $P_{B} = 4905 \text{ N/m}^2$
\n $P_{B} = 4905 \text{ N/m}^2$
\n $P_{B} = P_{B} = 4905 \text{ N/m}^2$
\n $P_{B} = P_{B} = 4905 \text{ N/m}^2$
\n $P_{B} = P_{C} + 3 \text{ kg}$ (1.13) = 4905 + 0.14 = 9810 m (1.3)
\n $P_{B} = P_{C} + 3 \text{ kg}$ (1.14) = 4905 + 0.14 = 9810 m (1.15)
\n $P_{B} = 16883 \text{ N/m}^2$
\n $P_{B} = 16888 \text{ N/m}^2$
\n $P_{B} = 138 \text{ kg}$, $P_{B} = 1905 \text{ kg}$
\n $P_{B} = 138 \text{ kg}$
\n $P_{B} = 130 \text{ kg}$ (1.15) = 4905 + 0.14 = 9810 \text{ m}^2
\n $P_{B} = 13 \text{ kg}$
\n $P_{B} = 138 \text{ kg}$
\n $P_{B} = 13$

$$
P_{0} = P_{P}
$$
\n-17200 + 0.7 x1810 (15-11.6) + 9810 (11.6-8)
\n= 1.6 x 9810 x h₂ + la₁...\n
$$
\therefore h_{2} = 2.64
$$
\n
$$
\therefore E_{G} = 8.1 + 2.64 = 10.64
$$
\n
$$
P_{C} = P_{D}
$$
\n
$$
= 17200 + 0.7 x 9810 (15-11.6) + 9810 (11.6-8) + 9810(4)
$$
\n
$$
= .13.6 x 9810 x h3 + 64
$$
\n
$$
\therefore h_{3} = 0.605
$$
\nH.W.1. As shown in Figure below, determine the pressure at point A.
\npoint A.
\n
$$
P_{0.6} = 12.164 R_{0}
$$
\n
$$
P_{0.6} = 12.164 R_{
$$

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3.4 Forces on Plane Surfaces (Panels)

This section explains how to represent hydrostatic pressure distributions on one face of a panel with a resultant force that passes through a point called the center of pressure. This information is relevant to applications such as dams, gates and water tanks.

Shatt Al-Arab University College/Department Civil Eng. Fluid Mechanics-I /second Year Dr. Jasim Mohsin 2/ Hydrostatic force on a Submerged Inclined Plane Surface liquid free Let a plane \equiv surface surface "A" with its centroid at $(sin_{\alpha}=\frac{h}{u})$ c.g. be plared h_c h_{P} at an angle (x) with respect $, dA$ to liquid free $*c-3$. $X \in P$ $surface.$ c.g.=center of
gravity c.p.=center of pressure Consider an elementary shaded strip of area dA f at a depth h. $F = \int P \cdot dA = \int A \cdot dA = \int A \cdot dA$ $F = Y sin \alpha$ y.dA - 0 Since; $\int y \cdot dA = \int_{c} A - Q$
 $\begin{pmatrix} \frac{1}{2} - \frac{1}{2} \frac{1}{2} \cdot dA \\ \frac{1}{2} - \frac{1}{2} \cdot dA \end{pmatrix}$ where; y =: centroid of plane surface a bout 0-0 axis $substitute$ eq. $@$ into eq. $@:$ \therefore F= δ $\Delta_c \cdot$ sin α . $A = \circledS$ $sinc_i$, y_c . $sinc_i$ h_c $\therefore \sqrt{F} = \sqrt{\frac{1}{2}h_{c}A}$ $\left(4\right)$

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 $F = \frac{1}{2}$ and h_y dr , f r re c N). where; he = vertical depth from centroid of the plane surface to the lighted free surface Cm). A: area of the inclined plane surface (a). Location of the Hydrastatic force: The point of action of the total hydrostadic force on the Surface is called the center of pressure (c.p.). This point is calculated by equating the moment of the total hydr. force acting at c.p. to the summation of the moments due to the elementary forces acting on the elementary strips; By applying Varignon's theorem; $CM_R = \sum_{\text{components}}$

 $F. y_P = JdF. y = Jp. dA. y = J8h. dA. y$: F. $y_p = \int Y y^2 sin\alpha \cdot dA = \delta sin\alpha \int y^2 dA$ $y_p = \frac{8 \sin \alpha \int_{A} y^2 dA}{\sqrt{2 \sin \alpha \int_{A} y^2 dA}}$ From $eq.0$: $F = 8 sin \alpha \int 9. dA$ $y_p = \frac{y_{sinx} \int_{A}^{y^2} dA}{y_{sinx} \int_{A}^{y} dA} \Rightarrow y_p = \frac{\int_{A}^{y^2} dA}{\int_{A}^{y} dA}$ Jy². dA = second moment of area of the plane surface about

 $0-0$ axis.

= moment of inertia of the plane surface about 0-0 axis.
\n=
$$
\Sigma_{0-0}
$$

\n \therefore Eq. (5) becomes : $yp = \frac{\Sigma_{0-0}}{y_{c/A}} - 6$
\nsince; $\Sigma_{0-0} = \Sigma_{c} + \frac{y_{c}^{2}}{f} \cdot A$
\nwhere; Σ_{c} : Mount of inertia of the plane surface
\nabout its centroid.
\n \therefore $y_{p} = \frac{\Sigma_{c} + y_{c}^{2} \cdot A}{y_{c/A}} \Rightarrow \boxed{y_{p} = \frac{\Sigma_{c}}{y_{c} \cdot A}}$

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3/ Hydrostatic force on a Submerged Curved Surface

On a curved surface, as shown below, the element force (dF) varies both in magnitude f in direction. The x f y- components of total force (F) can be evaluated by summation of elemental force components.

 \therefore Eq. \bigoplus becomes; $F_H = \delta h_c \cdot Av$

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subs. $=1.5$ into $=9.0$: $Fv = \int_{A} p. dA \cdot cos\theta = \int_{A_H} h. (dA)_A = \int_{H} 8 dH$ \therefore $Fu = \gamma +$ $F = \sqrt{F_H^2 + F_V^2}$ $F : Total hydr.$ force (N) . where; (dA): the projection of the element area normal $the x-axis$. (dA)H: the projection of the element area normal to the y-axis. Direction of the total hydr. force

$$
\tan\Theta = \frac{F_H}{F_V}
$$
\n
$$
\Theta = \tan^{-1}\left(\frac{F_H}{F_V}\right) \text{ with the}
$$
\n
$$
\text{Vert} \cdot \text{del}
$$

For the sketch shown;
\n
$$
F_1 = \frac{1}{2} S_1 \cdot h_1 \cdot h_1 (1) = \frac{1}{2} S h_1^2
$$

\n $F_2 = S_1 \cdot h_1 \cdot h_2 (1) = S_1 h_1 h_2$
\n $F_3 = \frac{1}{2} S_2 h_2 \cdot h_2 (1) = \frac{1}{2} S_2 h_2^2$
\n $F_{\text{total}} = F_1 + F_2 + F_3$
\nTo find the location;
\n S_3 applying Variyno n² - Theorem about point A
\n $F \cdot (3) = F_1(\frac{2}{3} h_1) + F_2(\frac{h_2}{2} + h_1) + F_3(\frac{2}{3} h_2 + h_1)$
\n $\therefore S = V$

Ex1: For the gate's shown in the figure below, calculate:

\n1- Hydrostate. Force on the gate's shown in the Figure below, calculate:

\n2- Turing moment about the axis of rotation.

\nSince
$$
F = 8 h_0 A
$$

\nY = Y_w = 9810 N/m³

\n= 9.81 K N/m³

\nFrom the sketch:

\nh = 2 - K

\n= 2 - $\frac{125}{2}$ sin80°

\n= 1.384 m

\nA = $\frac{x}{4}D = \frac{x}{4}$ (1.25)⁴

\nF = 9.81 * 1.384 * 1.227

\n= 16.66 KN

\nSince $Y = \frac{1}{2}e = \frac{\frac{x}{6}}{6}$

\nTherefore, $Y = \frac{1}{3}eA = \frac{\frac{x}{6}}{16}$

\nSince $Y = \frac{1}{3}eA = \frac{\frac{x}{6}}{16}$

\nThus, $\frac{1.329}{61080} = 0.07m$

\nIntning moment about the axis of the axis of the axis of the axis.

\nSo, $\frac{1.329}{61080} = \frac{1.25}{\frac{1}{36}} \sin 80^\circ = \frac{1.25}{\frac{1}{36}} \sin 80^\circ$

\nSo, $\frac{1.25}{\frac{1}{36}} = \frac{9.6}{6}$

\nSo, $\frac{1.25}{\frac{1}{36}} = \frac{9.6}{6}$

\nSo, $\frac{1.25}{\frac{1.2$

Ex.2 - The triangular gate CDE is hinged along CD and is opened by a normal force R applied at E. It holds oil (s. = 0.8) above it and is open to atmosphere on its lower side. The gate weighs 20kN. Find @ the magwithde of the hydrostatic force, (b) the location of pressure center, of @ the force R needed to open the gate. Solution. $\frac{6}{x}$. a) since F= 8 hc A $y = 0.8 * y_{\text{order}} = 0.8 * 9810$
= 7848 N/m³ ぬう $=7.84RKN/m^{3}$ $h_c = h_1 + 2$ $\overline{4}$
 $\overline{4}$
 $\overline{4}$ $h_c = \frac{2}{3}$ \star 5 \star sin 30 + 2 = 1.667+2 $= 3.667 m$ $A = \frac{1}{2}bh = \frac{1}{2}k3*5 = 7.5 m^2$ $5. F = 7.848 * 3.667 * 7.5 = 215.84$ kN $y_{P-J_c} = \frac{T_c}{y_{c.A}}$ $y_p = y_c + \frac{I_c}{y_c A}$ \bigodot $y_c = \frac{hc}{\sin 3a^2} = \frac{3.667}{\sin 3a^8}$ $sin 36 = \frac{hc}{y}$ \therefore $y_c = 7.334$ m

$$
L_{c} = \frac{bh^{3}}{36}
$$

\n
$$
\therefore L_{c} = \frac{3(5)^{3}}{36} = 10.417 \text{ m}^{4}
$$

\n
$$
\therefore \frac{L_{c}}{3cA} = \frac{10.417}{7.334*7.5} = 0.189 \text{ m}
$$

\n
$$
\therefore \frac{v_{c}}{9c} = 7.334 + 7.5 = 0.189 \text{ m}
$$

Shatt Al-Arab University College/Department Civil Eng. Fluid Mechanics-I /second Year Dr. Jasim Mohsin Ex.3. How long will the water on the right (h) has to rise to open the gate shown below. The gate is 2m wide, and is constructed of material with 3. = 4.5. w ata $1m$ Solution: $\overline{\mathbf{k}}$ For Fil By using press. dist. diagram $2m$ $F_1 = \frac{1}{2}$ (base) * (height) * b hinge $base = \delta_{w}(1) = 9.81kN$ $F_1 = \frac{1}{2} * 9.81 * 1 * 2 = 9.81$ KN Waate $\frac{3}{5}$ $\frac{2}{3}$ $\frac{1}{3}$ = 0.667m h $H.\omega$ use $F_1 = \gamma h c A_1$ F_1 $\frac{1}{\frac{1}{3}}$ $h = \frac{1}{2}$ $A = 1*2$ hrigh $F_1 = 9.81 + \frac{1}{2} \times 2 = 9.81 \text{ KN}$ $\overline{\mathcal{C}_{\omega}(1)}$ $y_{p=1} = \frac{2}{1} + \frac{2}{1} = 0.5 + \frac{2 \times 1^3}{0.5(1 \times 2)}$ $= 0.5 + \frac{2}{12} = 0.5 + 0.1667$ $For F2 : F5 = \int_{a}^{b} h_5 A_2$ $Cos\theta = \frac{1}{2}$ 700060 $F_2 = 9.81 * h_{c_2} * (2 * 2) = 39.24 h_{c_2}$ (KN) $sin60 = \frac{h}{1}$ $y = y_c + \frac{I_{c1}}{3cA} = 0$:. $h_1 = 0.886$ m $\therefore \frac{\Gamma c_2}{3c_2 A_2} = \frac{2(2)^3/12}{1.155 h_c (2 \times 2)} = \frac{0.288}{h c_2}$ $sin60^\circ = \frac{hc}{3c}$ $\therefore y_c = \frac{hc_2}{sin 60}$ 44 = 1.155 hc

$$
W = m \cdot 9; P = \frac{w}{4} \Rightarrow m = \rho +
$$

\n
$$
W = \rho 9 + \frac{1}{4} \times \frac{1}{
$$

$$
\sum M_{hinge} = 0.
$$
\n
$$
F_{2} * [1 - (3p_{2} - 3c_{2})] = F_{1} * \frac{1}{3} + \frac{W_{0}}{9} + \frac{1}{3}
$$
\n
$$
3q.24 h_{2} [1 - \frac{0.288}{h_{2}}] = \frac{9.81}{3} + \frac{76.46}{3}
$$
\n
$$
3q.24 h_{2} = 11.3 = 28.756
$$
\n
$$
\therefore h_{2} = 1.021 \text{ m}
$$
\n
$$
\text{since } h = h_{1} + h_{2}
$$
\n
$$
\therefore h_{2} = 0.886 + 1.021 = 1.89 \text{ m}
$$

$$
A = \int_{3}^{b} (1-y) dx
$$
\n
$$
A = \int_{3}^{b} (1-y) dx
$$
\n
$$
A = \int_{1-y}^{b} (1-y) dx = 0
$$
\n
$$
A = \int_{3}^{b} (1-x^{2}) dx = x \int_{0}^{1} -\frac{1}{3}x \int_{0}^{3} = 1 - (\frac{1}{3}) = \frac{2}{3}x
$$
\n
$$
\therefore \forall y = \frac{2}{3}x2 = \frac{4}{3}x
$$
\n
$$
\therefore \forall y = 9.81 \times \frac{4}{3} = 13.1 \text{ km } 4
$$
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$$
\frac{4}{3} = 13.1 \text{ km } 4
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\frac{4}{3} = 13.1 \text{ km } 4
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\frac{4}{3} = 13.1 \text{ km } 4
$$
\n
$$
\frac{4}{3} =
$$

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\nEx.s.

\nCalculate 1/secondYear

\nthe
$$
4\pi
$$
?

\nthe 4π

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Ex. 6: Find the Vertical component of force in the metal
\nspherical done shown in Figure below, when gauge A
\nreads 69 kfa. Assume the down weight 4500N
\n

\nNote. The volume of sphere =
$$
\frac{\pi D^3}{6}
$$

\nSubte. The volume of sphere = $\frac{\pi D^3}{6}$

\n3

\n3

\n3

\n3

\n6

\n4

\n5

\n5

\n6

\n6

\n7

\n8

\n1

\

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3.5 The Buoyant Force Equation Volume \ddot{r}_D The buoyant force (FB) passes through the center of buoyancy (B). $\frac{0}{1}$ Submerged Body Fuz 7 Fu_l because pressure increase F_{v} with depth
 $Fv_2 = 84z$ $*^{FB}$ $Fv_1 = 84$ $L_{\rm cl}$ where, δ = specific weight of the $\begin{picture}(120,20) \put(0,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}} \put(15,0){\line(1,0){155}}$ $H = H$ $H_{2} = H_{KLMNOK}$ t'2-t'1 = volume of displaced fluid = volume of submerged $body = #$ $E_{v_{2}} - F_{v_{1}} = 84 = F_{B}$ A^i $\cdot \sqrt{F_B} = \gamma + 1$ $\sqrt{\frac{F_{V}}{F_{V}}}$ $F_{\nu_{2}}$ $v_1 + v_2 -$ Ξ $t = t_2 - t_1$ $F_{B} = Y +$ $-49-$

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Stability of Immersed Bodies

When a body is completely immersed in a liquid, its stability depends on the relative positions of the center of gravity of the body and the centroid of the displaced volume of fluid, which is called the *center of buoyancy.*

- If the center of buoyancy is above the center of gravity (Fig. 3.11a), any tipping of the body produces a righting couple, and consequently, the body is stable.
- If the center of gravity is above the center of buoyancy (Fig. 3.11c), any tipping produces an increasing overturning moment, thus causing the body to turn through 180°.
- Finally, if the center of buoyancy and center of gravity are coincident, the body is neutrally stable—that is, it lacks a tendency for righting or for overturning, as shown in Fig. 3.11*b*.

Fig. 3.11

Stability Floating Bodies

The stability for floating bodies than for immersed bodies is very important because the center of buoyancy may take different positions with respect to the center of gravity, depending on the shape of the body and the position in which it is floating. When the center of gravity *G* is above the center of buoyancy *C* (center of displaced volume) for floating body, the body will be stable and equilibrium. The reason for the change in the center of buoyancy for the ship is that part of the original buoyant volume, as shown in Fig.3.12by the wedge shape *AOB*, is transferred to a new buoyant volume *EOD.* Because the buoyant center is at the centroid of the displaced volume, it follows that for this case the buoyant center must move laterally to the right. The point of intersection of the

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lines of action of the buoyant force before and after heel is called the *metacenter*

M, and the distance *GM* is called the *metacentric height.*

 \Box If *GM* is positive—that is, if *M* is above *G*, the body is stable \Box If *GM* is negative, the body is unstable.

Fig.3.12

Consider the prismatic body shown in Fig. 3.12, which has taken a small angle of heel α . First evaluate the lateral displacement of the center of buoyancy CC ^{*'*}, then it will be easy by simple trigonometry to solve for the metacentric height *GM* or to evaluate the righting moment.

The righting couple =W. MG. sin α

Where : *W* weight of body and α angle of heel.

By similar triangle *EOD* and *C'CM*: $\frac{\Delta y}{b/2} = \frac{C/C}{CM}$ find *CM*

 $GM = MC - GC$

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Ex.1	A spherical busy has adia. of (1.5m), weighs 8.5 KN, and is attached as shown in figure below with a cable. Determine the tension force at the cable. Note: volume of sphere = $\frac{\pi D^3}{6}$ where = $\frac{\pi D^3}{6}$ where = 14.80 of the begin
Suble	Example
Suble	Table
Suble	Step 1
Suble	Step 1
Suble	Step 25
Suble	Step 3
Suble	Step 4
Suble	Step 4
Suble	Step 5
Suble	Step 6

1- the deep it will sink
2. the mass of stone placed on the box to sink it
4m depth.

Ex3: An object weighs 3N in Water and 4N in oil
\n(s=0.83). Determine its volume f-specific gravity (s-).

\nSo.

\nSo.

\nSo.

\nW_{air} = W_{air} - Y_{wi} + ...

\nW_{irp}:1

\nW_{air} = W_{air} - Y_{ci}: + ...

\nW_{irp}:1

\nW_{air} = W_{air} - 1

\nW_{irp}:1

\nW_{air} = W_{air} - 1

\nW_{air}: 3 = W_{air} - 1

\nSubs. Eq. (3) into 77.0

\n4 = 3 + 1810 + ... = 23 (1810)(4)

\n4 = 3 + 1810 + 6 × 10

\n4 = 3 + 1810 + 6 × 10

\nNow Eq. (3)
$$
\Rightarrow
$$
 W_{air} = 3 + 1810 + 6 × 10

\nNow Eq. (3) \Rightarrow W_{air} = 3 + 1810 + 6 × 10

\nNow Eq. (4) \Rightarrow 6 × 10

\nNow Eq. (5) \Rightarrow W_{air} = 3 + 1810 + 6 × 10

\nNow Eq. (6) \Rightarrow 100/34

\n8.886 = $\begin{bmatrix} 6 & 6 & 16 \\ 6 & 6 & 16 \end{bmatrix}$

\nSo.

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Example 2.13 In Fig. a scow 20 ft wide and 60 ft long has a gross weight of 225 short tons (2000 lb). Its center of gravity is 1.0 ft above the water surface. Find the metacentric height and restoring couple when $\Delta y = 1.0$ ft.

SOLUTION

1. Find the depth (h):

The depth of submergence h in the water is

$$
h = \frac{225(2000)}{20(60)(62.4)} = 6.0
$$
 ft

2. Find the location of new center of buoyancy (C'):

The centroid in the tipped position is located with moments about AB and BC

$$
x = \frac{5(20)(10) + 2(20)(\frac{1}{2})(\frac{20}{3})}{6(20)} = 9.46 \text{ ft}
$$

$$
y = \frac{5(20)(\frac{5}{2}) + 2(20)(\frac{1}{2})(5\frac{2}{3})}{6(20)} = 3.03 \text{ ft}
$$

$$
y = \frac{1}{2}
$$

$$
y = 1
$$

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4.Find MG

 G is 7.0 ft from the bottom; hence,

$$
\overline{GC} = 7.00 - 3.03 = 3.97 \text{ ft}
$$

 $\overline{MG} = \overline{MC} - \overline{GC} = 5.40 - 3.97 = 1.43$ ft and

The scow is stable since \overline{MG} is positive; the righting moment is

$$
W\overline{MG}\sin\theta = 225(2000)(1.43)\frac{1}{\sqrt{101}} = 64,000 \text{ lb} \cdot \text{ft}
$$

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3.6 Equilibrium of accelerated fluid masses

If a body of fluid is moved at a constant velocity, then it obeys the equations derived earlier for static equilibrium.

If a body of fluid is accelerated such that, after some time, it has adjusted so that there are no shearing forces, there is no motion between fluid particles, and it moves as a solid block, then the pressure distribution within the fluid can be described by equations similar to those applying to static fluids.

Liquid in A Container Subjected to A Constant Rotation A liquid, contained in a vessel, may be rotated at a constant rotational velocity (w) without any relative movement being created between different elements of the liquid in the vessel. The liquid reprients itself once 4 for all to stay in that position with respect to the axis of rotation. Tangent 4° W new free Surface ω^2 . \times any
point \sqrt{v} $tan\theta =$ $w^2 \cdot x^2$ Parabola $y =$ Equation (new free surface $eg.$) wherei θ = Inclination angle of tangent of any point located along new free $axis of$ Surface (degree). $x, y = x$ $f(y) = x$ alues for any point located along new free surface. 68

69

$$
h_{max.} = 2 + 0.764 = 2.764m
$$
\n
$$
h_{min.} = 2 - 0.764 = 1.236m
$$
\n
$$
h_{min.} = 2 - 0.764 = 1.236m
$$
\n
$$
f_{1} = \frac{1}{2}h_{max.} = \frac{1810}{1000}(2.74)
$$
\n
$$
\therefore f_{1} = 27 \text{ KPa}
$$
\n
$$
f_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b = 149.2 \text{ KN } (+)
$$
\n
$$
P_{2} = \frac{V_{1}}{2}h_{max.} \cdot b = 149.2 \text{ KN } (+)
$$
\n
$$
P_{3} = \frac{12.125 \text{ KPa}}{2} \qquad \frac{100 \text{K}}{2} \cdot \frac{1}{4}h_{max.} \text{ k-fins.}
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{min.} \cdot b = 30 \text{ KN } (-+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{min.} \cdot b = 30 \text{ KN } (-+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{min.} \cdot b = 30 \text{ KN } (+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b = 30 \text{ KN } (+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b = 30 \text{ KN } (+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b = 30 \text{ KN } (-+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b = 30 \text{ KN } (-+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b = 30 \text{ KN } (-+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b = 30 \text{ KN } (-+)
$$
\n
$$
F_{1} = \frac{1}{2}h \cdot h_{max.} \cdot b
$$

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C =
$$
P_1 = P_2 = \gamma h (1 - \frac{a_4}{3}) = 13.62
$$
 KPa
\n $F_1 = F_2 = \frac{1}{2} P_1 (2) (4) = 54.48$ KN
\n $F_3 = P_1 (5) (4) = 272.4$ KN

 $d-$

$$
F_2 = 45.93 KN
$$

 $F_3 = \frac{P_1 + P_2}{2}(5)(4) = 452.4 KN$

Ex.2. If the tank in Ex.1 is accelerated horizontally along the longer side, determine the maximum acceleration that can be given without spilling the water. Also, calculate the percentage of water spilt if this max acceleration is increased by 20%.

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$$
\therefore \text{ The volume of water that Split over } = \frac{4}{\text{original}} - \frac{4}{\text{Keptin}}
$$

$$
\therefore \frac{9}{6} \text{ of water spin} + \text{over } = 4\frac{3}{\text{m}}
$$

$$
\therefore \frac{9}{6} \text{ of water spin} + \text{over } = \frac{4}{\text{m}} \times 100 = 10\%
$$

Ex.3: The tank in figure is filled with oil (s.=0.8) of acceleration as shown. There is a small opening in the tank at A. Determine the pressure at $B \not\subset C$; and the acceleration (ax) required to make the pressure at B equels (7 KPa "Vacuum").

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I surface $Sol.$ $\frac{y}{2} \frac{\frac{1}{4}}{4}$ $tan\theta = \frac{q_x}{q}$ Ve $9x = 4.903$ \leftrightarrow $\frac{1}{2}$ $\frac{ax}{9} = \frac{y_1}{1.8} = \frac{y_2}{9.15}$ new free
Surface $\mathbb B$ \therefore $y = 0.9$ m \overline{C} $-1.8 - +$ $y_{2} = 0.075m$ $l \rightarrow$ $0 - 15$ $P_B = \gamma \left(1.2 - 3\right) = 2.35 \text{ K}P_A$ $P_{c} = \delta_{c1} \left(1.2 + 0.15 + 9_{2} \right) = 11.18 K R.$ $-1.84 \blacklozenge$ when $P_{B}=-7KP_{A}$ A \int_{0}^{∞} convert PB to oil depth $qx=7$ -7 $*10^{3} = 8.1$ hoil 1.2_w new free
Surface(2) : $h_{ij} = -0.892m$ B \subset $tan\theta_2 = \frac{9x}{9} = \frac{1.2+0.892}{1.8}$ 0.892 : $9x = 11.4 \text{ m/s}^2$

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Ex. 4. An open cylindrical tank (2m) high f (1m) in diameter contains (1.5m) of water. If the cylinder rotates about its geometric axis, (a) hat constant angular velocity in (r.p.m) can be attained without spilling any water? (b) what is the pressure at the tank bottom (at C f D points) when $w = 6$ rad/s.? new free surface $\leq d$: (a) original \mathbf{y} Volume of paraboloid of revolution = $\frac{1}{2}$ (volume of circumscribed cylinder) 1.5 If no liquid is spilled, this $H - 41m$ volume equals the volume above the original free surface

Volume of paraboloid of revolution = $\frac{1}{2}(\overline{x}\overline{x}y)$ Volume above the original free surface = π (0.5) (0.5)

since there is no spilling of water
\n
$$
\frac{1}{2}(\pi (cos^{2}x)) = \pi (cos^{2}cos^{2}x)
$$
\n
$$
\therefore \frac{1}{2} = 1
$$
\n
$$
sincc : y = \frac{sin^{2}x^{2}}{29}
$$
\n
$$
\therefore w = 8.86 rad/s
$$
\n
$$
w = 8.86 rad/s
$$

From 4. ①
$$
\Rightarrow
$$
 $\mathbb{Z} = 0.23$ m
\n $\therefore P_{c} = \gamma_{w} (1.5 - (y - z)) = 12.46$ kPa
\n $P_{D} = \gamma_{w} (1.5 + z) = 16.97$ kPa
\nEx.S. Consider the tank in Ex.A closed with air space
\nsubjected to a pressure of (1.07 bar). When the
\nangular velocity is (12 rad. (s), what are the pressure
\nin (bar) at points C. 4 D?
\nSo.
\n \therefore From Ex.A, the w-value that
\nwalks the water reach the
\nbank bp edge is 8.86 rad./s.
\nSince $w = 12$ rad./s. ≥ 8.86 rad./s.
\n $\therefore w = 12$ rad./s. ≥ 8.86 rad./s.
\n $\therefore w = 12$ rad./s. ≥ 8.86 rad./s.
\n $\therefore w = 12$ rad./s. ≥ 8.86 rad./s.
\n $\therefore w = 12$ rad./s. ≥ 8.86 rad./s.

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$$
hp_i = y_i
$$

\nsince $y_i = \frac{y_i^2 + y_i^2}{2g_i^2}$
\n $= y_i = \frac{(|z_i|^2 (0.5))^2}{2g_i^2} = 1.83 \text{ m}$

 $3.92 = 1.31$ bar

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Chapter Four Fluid Dynamics

The force balance between pressure and weight in a static fluid was presented in unit 3, which led to an equation for pressure variation with depth. This unit deal with behaviour of fluid like velocity, acceleration and flow pattern (flow classification). The fluid dynamic deal with fluid having accelerated movement and there is relative acceleration between fluid particles.

Types of Fluid There are two types of fluid - Perfect (Ideal) (Non-Viscous)
There are two types of fluid - Real (viscous) fluid

 $-$ Perfect fluid \Rightarrow Viscosity =0. x In perfect fluid there is
a slipped Boundary condition

 $-$ Real $f(u, d) \Rightarrow V$ 'scasity $\neq o$. * In Real fluid, there is no slip Boundary condition due to viscous

Types of Flow:

\n1. Steady Flow:

\nexists if the velocity at a point, for example, remaining constant with respect to time
$$
(\frac{dy}{dt} = 0)
$$
.

\n2. Unsteady Flow:

\nexists if the velocity at a point, for example, changes either in magnitude or in direction with respect to time $(\frac{dy}{dt} + 0)$.

\n3. Uniform Flow:

\nexists if the velocity, for example, remains constant with respect to distance $(\frac{dy}{ds} = 0)$.

\n4. Non- Uniform Flow

\n1. The velocity of the velocity is the change of the distance of the velocity.

\n2. Using the velocity of the velocity, for example, changes either in magnitude or in direction, with respect to distance $(\frac{dy}{ds} = 0)$.

\n3. Using the velocity of the velocity, the example, changes either in magnitude or in direction, with respect to distance $(\frac{dy}{ds} = 0)$.

\n4. Non- Uniform Flow

\n5. The velocity of the velocity is the distance of the velocity.

\n6. The velocity of the velocity is the velocity of the velocity.

\n7. The velocity of the velocity is the velocity of the velocity.

\n8. The velocity of the velocity is the velocity of the velocity.

\n9. The velocity of the velocity is the velocity of the velocity.

\n1. The velocity of the velocity is the velocity of the velocity.

\n1. The velocity of the velocity is the velocity of the velocity.

\n1. The velocity of the velocity is the velocity of the velocity.

\n2. Unsteady Flow:

\n1. The velocity of the velocity is the velocity of the velocity.

\n2. The velocity of the velocity is the velocity of the velocity.

\n3. The velocity of the velocity is the velocity of the velocity.

\n4. The velocity of the velocity is the velocity of the velocity.

\n1. The velocity of the velocity is the velocity of the velocity.

\n2. The velocity of the velocity is the velocity of the velocity.

\n3. The velocity of the velocity is the velocity of the velocity.

\n4. The velocity of the velocity is the velocity of the velocity.

\n5. The velocity of the velocity is the velocity of the velocity.

\n1. The velocity of the velocity is the velocity of the velocity.

\n2. The velocity of the velocity is the velocity of the velocity.

\n3. The velocity of the velocity is the velocity of the velocity.

\n4. The velocity of the velocity is the velocity of the velocity.

\n5. The velocity of the velocity is the velocity of the velocity.

\n1. The velocity of the velocity is the velocity of the velocity.

\n2. The velocity of the velocity is the velocity of the velocity.

\n

Shatt Al-Arab University College/Department Civil Eng. Fluid Mechanics-I /second Year Dr. Jasim Mohsin Strain Line: Is an imaginary line within the flow for which the tangent at any point is the time average of the direction of motion at that point. streamline Stream Tube: Is an element of fluid bounded by a special group of streamlines which enclose or confine stream the flow. $f_{\mu}b$ e stream lines One, Two & Three Dimensional Flow - 1D-Flow: Such as flow in pipe 1D & Arisymmetry flow f_{low} 2D-Flow: Such as flow around the wing of aircraft. air flow 82

Velocity 4 Acceleration
\nMotion of fluid is specified by velocity components
\nexpressed as functions of space 4 times
\n
$$
u = F(x, y, z, t) - velocity component in x-dir.\n
$$
u = F(x, y, z, t) - v = v \qquad (x - 4i\pi)
$$
\n
$$
u = F(x, y, z, t) - v \qquad (x - 4i\pi)
$$
\n
$$
u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
$$
\n
$$
u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
$$
\n
$$
u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
$$
\n
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
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\n
$$
u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
$$
\n
$$
u = F(x, y, z, t) - v \qquad (x - 2 - 4i\pi)
$$
\n
$$
u =
$$
$$

where: Q = flowrate (discharge) in
$$
(2^{3}/T)
$$
 Letcept hand
\n $A = cross-sectional area in (L^{2})$
\n $V = average velocity in (L/T)$
\n Q_{1}, Q_{2}, LQ_{3} : How rate at sections (0, 2), 4 (3),
\nrespectively.

Solution:
\nSince,
$$
W = \gamma\psi
$$

\n $\therefore \gamma = \frac{W}{\gamma} = \frac{3000}{9810}$
\n $\therefore \gamma = 306 \text{ w}^3$
\n $\varphi = \frac{V}{t}$; for $+ = \{s, \Rightarrow Q = \frac{0.306}{1} = 0.306 \text{ w}^3/s.\}$
\n $\therefore Q = A \cdot V \Rightarrow V_1 = \frac{Q}{A_1} = \frac{0.306}{\frac{1}{4}(0.3)^2} = 4.33 \text{ w/s}$
\nSimilarly; $V_2 = \frac{Q}{A_2} = 9.74 \text{ m/s}$

Ex.2, As shown in figure below, if
$$
D_A = 450 \text{ mm}
$$
,

\n
$$
D_B = 300 \text{ mm}, D_C = 150 \text{ mm}, D_D = 225 \text{ mm}, \sqrt{A} = 1.8 \text{ m/s.}
$$
\n
$$
A = \sqrt{D} = 3.6 \text{ m/s.}
$$
\n
$$
A = \sqrt{D} = 3.6 \text{ m/s.}
$$
\nAdi>200

\nEquation:

\n
$$
A = \sqrt{A} = \sqrt{B} = \sqrt{C} + \sqrt{C}
$$
\n
$$
A = \sqrt{A} = \sqrt{B} = \sqrt{C} + \sqrt{C} + \sqrt{D} = \sqrt{D}
$$
\n
$$
A = \sqrt{A} = \sqrt{A} = \sqrt{B} = \sqrt{B} = \sqrt{C} = \sqrt{A} = \sqrt{D}
$$
\nFrom eq. ①: $A_A \cdot V_A = A_B \cdot V_B \Rightarrow \frac{\pi}{4} (0.45)(1.8) = \frac{\pi}{4} (0.35)^3 V_B$

\n
$$
\frac{\pi}{4} (0.45)^3 (1.8) = \frac{\pi}{4} (0.15)^3 V_C + \frac{\pi}{4} (0.225)^3 (3.6)
$$
\n
$$
V_C = 8.09 \text{ m/s.}
$$

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So, Power of pressure = $\delta Q \frac{P}{\delta} = QP$ 1, 4 elevation = $8Qz$ $\frac{v}{v}$ $\frac{v}{z} = \frac{\sqrt{Q}}{2} = \frac{\sqrt{2}}{2}$ 4 4 Pump = $8QHa$ $\frac{1}{4}$ Aurbine = δQH dissipation power due to friction = 00hL 88

For the Venturi meter shown in figur below, the $Ex.1$ deflection of the mercury in the differential gauge is 0.36m. Determine the flow of water through the meter if no energy is lost between A & B. $4F10W$ Sol.: Since there is no energy loss between A & B, by applying Bern. Eq. between points A f B: \mathbf{B} $\frac{4}{150}$ \bar{x} 0.75 $\frac{P_A}{X_1} + Z_A + \frac{V_A^2}{z_9} = \frac{P_B}{\gamma} + Z_B + \frac{V_B^2}{z_9}$ AX $\frac{4}{300}$ Take datum at $A \Rightarrow Z_A = 0$. $ZR = 0.75$ \overline{C} So, Bern. Eq. becomes Hq $\frac{P_A}{X_B} + \frac{V_A^2}{39} = \frac{P_B}{X_W} + 0.75 + \frac{V_B^2}{39}$ - 0 From the differential gauge; Pc = PD $PA + 8wX + 8w(0.36) = PB + 8w(0.75) + 8w \cdot X + 13.68w(0.36)$ $-(2)$ Dividing Eq. 2 by Yw: : $\frac{PA}{X_{\text{H}}} = \frac{P_{B}}{X_{\text{H}}} + 5.286$ - $\binom{3}{}$

From continuity =
$$
\frac{OA = QB}{AA \cdot VA = AB \cdot PB}
$$

\n $\therefore VA = \frac{AB}{AA}VB$
\n $\therefore VA = \frac{AB}{AA}VB$
\n $\therefore VA = 0.25VB$ (4)
\nSubs. eqs. (3) f (4) in the 9.0:
\n $\frac{PB}{Xw} + s.286 + \frac{(0.25V_B)^2}{2g} = \frac{PB}{Xw} + 0.45 + \frac{V_B^2}{2g}$
\n $\therefore VB = 9.74 \text{ m/s}$
\n $\Leftrightarrow QB = 0.17 \text{ m}^3/s$.

- i- the outlet flow. i- the pressures at points C, D, E, 4 F. iii- plot the E.L. of H.G.L.
- Note: Assume no energy loss.

Solve:
$$
1 - \frac{1}{2}
$$
 since there is no energy loss,
\nby applying Bern. Eq. between points
\n $(\theta, \frac{1}{2}, \frac{\pi}{6})$:
\n $\frac{P_{1}x}{\theta} + Z_{1} + \frac{V_{1}x}{\theta} = \frac{P_{1}x}{\theta} + Z_{1} + \frac{P_{1}x}{\theta} = \frac{P_{1}x}{\theta} + Z_{1} + \frac{P_{1}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{2} + \frac{P_{2}x}{\theta} = \frac{P_{1}x}{\theta} + Z_{1} + \frac{V_{1}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1} + \frac{V_{1}x}{\theta} = \frac{P_{1}x}{\theta} + Z_{1} + \frac{V_{1}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{1}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{1}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1} + \frac{V_{2}x}{\theta} = \frac{P_{2}x}{\theta} + Z_{1$

 $\frac{p_1}{8}$ +0+ $\frac{v_1^2}{29}$ +Ha = $\frac{p_2}{8}$ + $\frac{z_2}{29}$

 Hq

Since; Power =
$$
\frac{1}{2}
$$
 Q Ha
\n
$$
12*746 = 9810
$$
 Q rHa
\n
$$
4a = \frac{0.912}{Q} - Q
$$
\n
$$
3aQ - Q
$$
\n
$$
9a = Q_2 = Q \Rightarrow V_1 = \frac{Q}{A_1} = 31.83Q - 3
$$
\n
$$
V_2 = \frac{Q}{A_2} = 127.32Q - 4
$$
\nFrom the manometer's $P_c = P_0$
\n
$$
P_1 + V_{H_2}(0.89) = P_2 + V_0 Z_2 + V_0 (0.89) \div V_0
$$
\n
$$
= \frac{P_2}{V_0} + Z_2 = \frac{P_1}{V_0} + 11.214 - 5
$$
\nSubs. eqs. (2), (3), (4), 45 into eq. 0:
\n
$$
\frac{P_1}{V_0} + \frac{(31.83Q^2)}{29} + \frac{0.912}{Q} = \frac{P_1}{V_0} + 11.214 + \frac{(127.32Q^2)}{29}
$$
\n
$$
= \frac{71}{60} + 11.214Q - 0.912 = 0.
$$
\nBy trial 4 error \Rightarrow Q = 0.0814 m/s.